# Dark Matter Intro

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#### Dark Matter

- How do we know it exists?
- How is it distributed?
- How does it interact?
- What is it?
- How was dark matter produced?

#### How do we know it exists?

• Dark matter is inferred from its gravitational effects on galactic and cosmological scales.

**Observations:** 

$$G_{\mu\nu} \neq 8\pi G_N T_{\mu\nu}$$

 $T_{\mu
u}$ 

Dark Matter Dark Energy

 $G_{\mu\nu}$ 

Modify gravity

#### How do we know it exists?

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T<sub>μν</sub> Dark Matter Standard Dark Energy Cosmological Model

$$G_{\mu\nu}$$

Modify gravity

### A parable: Neptune and Vulcan





- In the I800's the perihelion of Uranus' orbit was not consistent with the known objects in the solar system.
- Dark Matter: Adams and Le Verrier postulated a new planet, Neptune, to explain the orbit of Uranus — it was found where predicted just months later!
- Modified Gravity: Building on this success, Le Verrier postulated a new planet, Vulcan, to explain the precession of Mercury's perihelion. Wrong! The correct explanation is GR.

#### A parable: Neptune and Vulcan





#### Keep an open mind!

#### An example: velocity dispersion

• Given a collection of point masses in orbit around their centre of mass, we can apply the Virial Th'm to obtain the total mass:



$$\langle KE \rangle = -\frac{1}{2} \langle PE \rangle$$
$$\frac{1}{2} m \langle v^2 \rangle \simeq \frac{1}{4} \frac{G_N Mm}{R}$$
$$M = \frac{2R \langle v^2 \rangle}{G_N}$$

#### Aside: astronomical distances



Universe.....

#### An example: velocity dispersion

$$M = \frac{2R\langle v^2 \rangle}{G_N}$$
$$G_N = 4.32 \times 10^{-3} \; (\text{km/s})^2 \text{pc/M}_{\odot}$$

#### **Globular Clusters**

 $R \sim 10 ~{
m pc}$  $\langle v^2 \rangle \sim 10 ~({
m km/s})^2$  $M \sim 10^4 ~{
m M}_{\odot}$ Compare with Luminosity:  $L \sim 10^4 ~{
m L}_{\odot}$ 

Can account for mass from stars!

#### **Galaxy Clusters**

 $R \sim 2 \,\,\mathrm{Mpc}$  $\langle v^2 \rangle \sim 10^6 \,\,(\mathrm{km/s})^2$  $M \sim 10^{15} \,\,\mathrm{M_{\odot}}$ Compare with Luminosity:  $L \sim 10^{13} \,\,\mathrm{L_{\odot}}$ 

Can't account for mass from stars!

### An example: velocity dispersion

- A solution: There is ~100 times more non-luminous matter (dark matter) inhabiting galaxy clusters than luminous matter. If this is distributed uniformly inside the cluster, then the amount inside globular clusters is negligible.
- Another solution: Gravity behaves differently in clusters of galaxies than clusters of stars.

#### Evidence for dark matter

• There is evidence for dark matter across a wide variety of scales:





dwarf galaxies  $M \sim 10^9 \,\mathrm{M_{\odot}}$ 

galaxies  $M \sim 10^{12} \, {\rm M_{\odot}}$  $R \sim \text{kpc}$   $R \sim 10 - 100 \text{ kpc}$   $R \sim 1 - 10 \text{ Mpc}$ 

galaxy clusters  $M \sim 10^{15} \, {\rm M_{\odot}}$ 



observable Universe

 $R \sim 100 - 10^4 \text{ Mpc}$ 

astrophysics: non-linear

cosmology: linear

### Modelling the Universe

- Much of cosmology is described by General Relativity and Relativistic Hydrodynamics of perfect fluids.
- The fundamental variables:

 $g_{\mu\nu}(x,t)$   $\rho_i(x,t)$   $u^{\mu}(x,t)$ metric fluid densities fluid velocities

• The laws: Einstein and continuity equations

 $G_{\mu\nu} = 8\pi G T_{\mu\nu} \qquad \nabla_{\mu} T^{\mu\nu} = 0 \qquad T_{\mu\nu} = \sum_{i} \left(\rho_{i} + p_{i}\right) u_{i\mu} u_{i\nu} + p_{i} g_{\mu\nu}$ 

Set of coupled PDE's -- need initial conditions!

#### The Linear Universe

• For much of the history of the Universe:

$$g_{\mu\nu}(x,t) = \bar{g}_{\mu\nu}(t) + \delta g_{\mu\nu}(x,t)$$
  

$$\rho_i(x,t) = \bar{\rho}_i(t) + \delta \rho_i(x,t) \qquad \delta = \text{ small}$$
  

$$u^{\mu}(x,t) = \bar{u}^{\mu} + \delta u^{\mu}(x,t)$$

- Why was the universe so nearly homogeneous?
- For today: this is an extraordinary convenience!



#### The Linear Universe

• The homogeneous Universe:

$$p_i = w_i \rho_i$$
  $ar{
ho}_i(t_0)$   
Types of fluids Densities today

• The Linear Universe:

Characterize statistics of inhomogeneities!

 $\mathcal{P}\left[g_{\mu\nu}\left(x,t=0\right)\right] \quad \mathcal{P}\left[\rho_{i}\left(x,t=0\right)\right] \quad \mathcal{P}\left[u_{i\mu}\left(x,t=0\right)\right]$ 

Assume our Universe is typical.

### 6 Parameter Model of the Universe





2 parameters

 $\mathcal{P}\left[\delta g_{\mu\nu}\left(x,t=0\right)\right]$ 

I parameter (for CMB)



 $\mathcal{T}$ 

### Modelling the universe

Initially small fluctuations collapse to form galaxies, stars, etc.



#### That's it!

The rest is details.

# Giving Thanks

- The non-linear Universe
  - GR is highly non-linear inferring the state of the early universe would be like asking for the weather 100 million years ago based on the weather today.
  - No general classification of metrics how to characterize initial conditions?
  - Shock waves, singularities, baryonic feedback, oh my!

## Giving Thanks

- The linear Universe
  - Simple evolution allows initial conditions to be inferred.
  - Background evolution and growth of structure can be analyzed separately.
  - Simple classification of initial conditions and metric degrees of freedom.
  - Physics on different scales evolves independently (Fourier modes independent).

#### The rest

#### Now for some details....



• The metric in a flat, homogeneous, isotropic universe:



constant comoving distance = growing physical distance

• Conformal time:  $\eta = \int \frac{dt}{a(t)}$   $ds^2 = a^2(\eta) \left[ -d\eta^2 + \delta_{ij} dx^i dx^j \right]$ 



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#### particle horizon



• Equations of motion in a homogeneous universe:

$$G_{\mu\nu} = 8\pi G T_{\mu\nu} \quad \square \searrow \quad H^2 \equiv \left(\frac{\dot{a}}{a}\right) = \frac{8\pi G\rho}{3}$$

$$\nabla_{\mu}T^{\mu\nu} = 0 \qquad \Longrightarrow \qquad \dot{\rho} = -3H\left(\rho + p\right)$$

 $p = w\rho$ 

$$\rho = \rho_0 a^{-3(1+w)} \qquad a(t) = a_0 t^{\frac{2}{3(1+w)}}$$

different fluids gravitate differently!



• Evolution of the scale factor:



• Energy budget:

$$\left(\frac{H}{H_0}\right)^2 = \sum_i \Omega_i a^{-3(1+w_i)} \qquad \sum_i \Omega_i = 1$$

 $H_0 = 67.4 \ \frac{\mathrm{km}}{\mathrm{s Mpc}}$ 

 $c/H_0 = 4450 \text{ Mpc}$ 



• For non-relativistic matter in flat space:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0 \qquad \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} + \frac{\nabla p}{\rho} + \nabla \phi = 0$$
continuity
Euler
$$\nabla^2 \phi = 4\pi G\rho$$
Poisson
Linearize:

$$\begin{split} \rho &= \bar{\rho} + \delta \rho \qquad \phi = \phi + \delta \phi \qquad \vec{v} = \bar{\vec{v}} + \delta \vec{v} \\ p &= \bar{p} + c_s^2 \delta \rho \\ c_s^2 &= \frac{\partial p}{\partial \rho} = w \end{split}$$

• Linearized equation of motion:

$$\frac{\partial^2 \delta \rho}{\partial t^2} - c_s^2 \nabla^2 \delta \rho - 4\pi G \bar{\rho} \delta \rho = 0$$

• Fourier transform:

$$\delta\rho(t,\vec{x}) = \int \frac{d^3k}{(2\pi)^3} \delta\rho(t,k) e^{i\vec{k}\cdot\vec{x}} \qquad \frac{\partial^2\delta\rho}{\partial t^2} + \left(c_s^2k^2 - 4\pi G\bar{\rho}\right)\delta\rho = 0$$

$$\delta\rho(t,k) = A\exp\left(i\omega(k)t\right) + B\exp\left(-i\omega(k)t\right)$$

$$\omega(k) = \sqrt{k^2 c_s^2 - 4\pi G\bar{\rho}}$$

$$\omega(k) = \sqrt{k^2 c_s^2 - 4\pi G \bar{\rho}}$$

• Jeans scale: competition between pressure and gravity

$$\lambda_J = \frac{2\pi}{k_J} = c_s \left(\frac{\pi}{G\bar{\rho}}\right)^{1/2}$$

$$\omega(k) = \sqrt{k^2 c_s^2 - 4\pi G \bar{\rho}}$$

• Jeans scale: competition between pressure and gravity



• In an expanding universe waves are stretched:



- Only gravitationally bound (non-linear) structures separate from the Hubble flow.
- Expansion inhibits collapse.

 $t_{\rm coll} \sim (4\pi G \bar{\rho})^{-1/2} \sim H^{-1}$  Expansion will be relevant!

• Including expansion (on small scales):

$$\begin{bmatrix} \frac{d^2}{dt^2} + 2H\frac{d}{dt} + \left(c_s^2 \frac{k_{\rm com}^2}{a^2} - 4\pi G\bar{\rho}\right) \end{bmatrix} \frac{\delta\rho(t, k_{\rm com})}{\bar{\rho}(t)} = 0$$

$$c_s^2 = 0$$

$$\frac{\text{radiation}}{\frac{\delta\rho}{\bar{\rho}} \propto \log(a)} \qquad \frac{\delta\rho}{\bar{\rho}} \propto a \qquad \frac{\delta\rho}{\bar{\rho}} \propto \text{const.}$$

#### Growth

- On very large scales, there is no growth need GR.
- Between  $t_{\rm eq} < t < t_0$ , on small scales, density can grow by a factor of about  $1/a_{\rm eq} \sim 3000$

$$\frac{\delta\rho}{\bar{\rho}}(t_{\rm eq}) > a_{\rm eq} = \frac{\Omega_r}{\Omega_m}$$

 If there were no dark matter, matter-radiation equality happens later and growth during matter domination would be a factor of 5 smaller!

#### The most important plot

• An important scale: comoving horizon

• Horizon crossing: k = aHcomoving scale horizon crossing  $\eta$ superhorizon subhorizon k

conformal time  $\,\eta\,$ 

1

 $\overline{aH}$ 

#### The most important plot

• An important scale: comoving horizon

• Horizon crossing: k = aH



conformal time  $\,\eta\,$ 

1

 $\overline{aH}$
- An important scale: comoving horizon
- Horizon crossing: k = aHcomoving scale  $\eta$ smaller scales collapse first

conformal time  $\,\eta\,$ 

 $\overline{aH}$ 

• An important scale: comoving horizon



conformal time  $\,\eta\,$ 

1

 $\overline{aH}$ 



• Horizon crossing: k = aH



conformal time  $\,\eta\,$ 

• An important scale: comoving horizon

• Horizon crossing: k = aH



conformal time  $\,\eta\,$ 

1

 $\overline{aH}$ 

## Initial Conditions

- The initial fluctuations are imprinted in the gravitational potential.
- Empirically, the potential is a random Gaussian field whose power spectrum (all the info!) is nearly scale invariant (equal power in each decade of scale):



$$\langle \phi(\vec{k})\phi(\vec{k}')\rangle = \delta^{(3)}(\vec{k} - \vec{k}')P_{\phi}(k)$$
$$P_{\phi}(k) = \frac{A_s}{k^3} \left(\frac{k}{k_*}\right)^{n_s - 1}$$

Adiabatic: density perturbations follow potential perturbations.

#### Inflation: the prequel to the Observable Universe



#### Matter Power Spectrum

Relating the potential fluctuations to density perturbations:



#### Baryons and photons

equ	ality recomb	ination reioni	zation
	photons tightly coupled to baryons imperfect fluid with $c_s^2\simeq 1/2$	neutral atoms form photons and baryons decouple: CMB is released!	first non-linear structure stars reionize the Universe
	dark matter perturbations begin to grow	baryons begin to collapse into dark matter halos	hierarchical structure formation some CMB photons re- scatter

 Before recombination, photons and baryons are coupled — pressure!



dark matter potential well

 Before recombination, photons and baryons are coupled — pressure!



 Before recombination, photons and baryons are coupled — pressure!



 Before recombination, photons and baryons are coupled — pressure!



 Before recombination, photons and baryons are coupled — pressure!



 After recombination, baryons decouple — scale of the sound horizon (distance sound travels since the big bang) is imprinted!

• Matter power spectrum:



Imprints in the distribution of photons (CMB):



#### Cold dark matter

• The model for linear growth of inhomogeneities described above assumed dark matter is composed of non-relativistic, non-interacting particles.

 $c_s = 0 \quad \longleftrightarrow \quad \text{cold dark matter (CDM)}$ 

- Massless dark matter particles would be 'hot' and behave like radiation — pressure! Not much growth and no acoustic oscillations, power suppression on scales beyond the peak of the matter power spectrum.
- Warm dark matter dark matter particles with small mass become non-relativistic some time during cosmic history. Power suppression on small scales, can preserve BAO.

#### Dark Matter

- Cosmology provides ample evidence for dark matter:
  - Luminous matter not enough to account for growth.
  - BAO peaks in matter power spectrum and CMB angular power spectrum require dark matter.
  - Much more! For example, weak lensing of CMB.





#### Nonlinear scales



- Inhomogeneities smaller than ~ 50-100 Mpc are nonlinear.
- These are the 'structures' that fill the Universe; in fact, it is not a bad approximation to assume that all dark matter is bound up into halos of varying mass.

#### Formation of dark matter halos

- Primordial dark matter perturbations are a random Gaussian field.
- Growth will drive peaks in the distribution into the nonlinear regime — gravitational collapse happens in these regions.



#### Formation of dark matter halos

- Three stages of collapse:
  - Linear collapse: ~linear growth of overdensity to critical value of  $~\delta
    ho/\bar
    ho\simeq 1.686~$ .
  - Virialization: non-linear evolution to a virialized halo.
  - Secondary infall: surrounding dark matter is accreted.

#### **NFW Profile**

• The result: NFW profile



• A reasonable fit to N-body simulations.

#### **Rotation Curves**

• Circular orbital velocity:

$$V_c(R) = \sqrt{\frac{G_N m(R)}{R}}$$



#### Halo Mass Function

• The number density of halos as a function of mass:



#### Halo Bias

Halos are more likely to form on the peaks of the DM distribution:



 Halos in a fixed mass range don't form at random positions, but rather trace the underlying dark matter distribution.

$$P_h(k|m) = b_h(m)^2 P_m(k)$$

#### Halo Model of LSS

• If we assume all dark matter is bound inside halos:



## Adding Galaxies

 Halo Occupation Distribution: prescription for adding galaxies to halos; assumed to depend primarily on halo mass.



## Small-Scale issues with LCDM

- Core cusp problem: galaxy rotation curves of low-mass galaxies prefer a constant central dark matter density.
- Diversity problem: a larger than expected scatter in galaxy central densities inferred from rotation curves.
- Missing satellites problem: far fewer small galaxies in local group are observed than expected from the number of low-mass dark matter halos.
- Too-big-to-fail problem: Dwarf galaxies should have more stars than observed, based on their dark matter density.

#### **Baryonic effects or new physics?**

#### How does dark matter interact?

• Beyond gravity, does dark matter interact with itself or with baryonic matter?

DM is motivated by gravitational effects. Why consider anything else?

- Compelling theory: dark matter is part of a more complete theoretical construction that necessarily predicts additional interactions.
- Empirical: observations require additional interactions.
- Nature is kind: without additional interactions, it will be difficult to determine what dark matter is.

## Compelling theory?

- Best example: supersymmetry.
- Solves hierarchy problem, but includes weakly interacting massive particles (WIMPS) that can be the dark matter.

#### But SUSY in trouble....

- QCD Axion: can solve strong CP problem and serve as a dark matter candidate.
- Sterile neutrino: can be part of the mechanism responsible for neutrino masses.

In these scenarios, dark matter can interact with baryons and/or itself.

#### **Empirical evidence?**

• Some of the small-scale problems with LCDM can be alleviated with self-interacting dark matter.

#### Could be baryonic physics.....

• Limits on interactions:









merging clusters

direct detection

indirect detection

direct production

And many more creative ways....

#### Nature is kind?

- I really hope that dark matter has non-gravitational interactions!
- If nature is kind: what interactions can we test given existing technology?

Can drive a creative cycle: we won't know how to look if we don't know what to look for.

Can be dangerous: if dark matter could be anything, and interact in any way, where do we stop?

#### What is dark matter?

- Some possibilities:
  - Fundamental particle(s).
  - Composite particle(s).
  - Condensate(s).
  - Black holes.
  - •

Important and interesting to look for observations that can distinguish these various candidates.

#### How was dark matter produced?

- A few possibilities:
  - Thermal production (e.g. WIMP)

• Non-thermal production (e.g. Axion)

#### Thermal production



$$\Omega_{\chi} = 0.7 \left( \frac{m_{\chi}/I_f}{10} \right) \left( \frac{g_*(m_{\chi})}{100} \right)^{-1} \left( \frac{10^{-5c} \text{ cm}^2}{\langle \sigma v \rangle} \right)$$

#### **WIMP Miracle:**

 $\Omega_{\chi} \sim 0.3 \quad m_{\chi} \sim \text{TeV} \quad \langle \sigma v \rangle \sim 10^{-39} \text{ cm}^2$ 

#### Non-thermal production



 $\mathcal{A}$ 

- When m > H, a oscillates and dilutes like pressureless dust.
- Gravity and/or self interactions cause the condensate to fragment.

#### Some perspectives/opinions





# Should we modify gravity?



Is dark matter one component, or is there a dark sector?

#### Some perspectives/opinions



# How will we identify what dark matter is?